

Chapter 21

Integral Equations

Introduction.

linear.

$$f(x) = \int_a^b K(x, t) \psi(t) dt \quad |$$

Fredholm equation

$$\psi(x) = f(x) + \lambda \int_a^b K(x, t) \psi(t) dt \quad 2$$

$$f(x) = \int_a^x K(x, t) \psi(t) dt \quad |$$

Volterra equation

$$\psi(x) = f(x) + \int_a^x K(x, t) \psi(t) dt \quad 2$$

Transformation:

Differential \rightarrow Integral

Linear second-order ordinary differential equation.

$$y'' + A(x)y' + B(x)y = g(x)$$

initial conditions:

$$y(a) = y_0 \quad y'(a) = y'_0$$

Integrate.

$$\begin{aligned} y'(x) &= - \underbrace{\int_a^x A(t)y'(t) dt}_{-A(x)y(x)} - \int_a^x B(t)y(t) dt + \int_a^x g(t) dt + y'_0 \\ &= -A(x)y(x) - \int_a^x [B(t) - A(t)]y(t) dt + \dots \end{aligned}$$

Integrate again.

$$y(x) = - \int_a^x A(t)y(t)dt - \int_a^x du \int_a^u [B(t) - A'(t)] y(t) dt \\ + \int_a^x du \int_a^u g(t) dt + (A(a)y_0 + y'_0)(x-a) + y_0$$

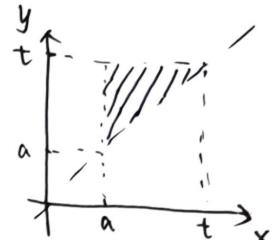
To make it neater,

$$\int_a^x du \int_a^u f(t) dt \rightarrow \int_a^t dy \int_a^y f(x) dx$$

$$\int_a^x f(t) dt \int_t^x du \leftarrow \int_a^t dx \int_x^t f(x) dy$$

"

$$\int_a^x f(t) \cdot (x-t) dt$$



$$\text{So, } y(x) = - \underbrace{\int_a^x (A(t) + (x-t)[B(t) - A'(t)]) y(t) dt}_{\rightarrow K(x,t)} \\ + \underbrace{\int_a^x (x-t) g(t) dt + [A(a)y_0 + y'_0](x-a) + y_0}_{\rightarrow f(x)}$$

$$\Rightarrow y(x) = f(x) + \int_a^x K(x,t) y(t) dt$$

Volterra equation of the second kind.

example : linear oscillator

$$y'' + \omega^2 y = 0$$

$$y(0) = 0, \quad y'(0) = 1$$

$$\Rightarrow A(x) = 0 \quad B(x) = \omega^2 \quad g(x) = 0$$

So

$$y(x) = x + \omega^2 \int_0^x (t-x) y(t) dt$$

It's easy to check that

$$y(x) = \frac{1}{\omega} \sin(\omega x)$$

is its solution.

If the boundary condition becomes

$$y(0) = 0 \quad y(b) = 0$$

We don't know $y'(0)$, so, Integrate.

$$y' = -\omega^2 \int_0^x y dx + y'(0)$$

Integrate again.

$$y = -\omega^2 \int_0^x (x-t) y(t) dt + x y'(0)$$

use $y(b) = 0$,

$$\omega \int_0^b (b-t) y(t) dt = b y'(0)$$

Substitute back.

$$y = -\omega^2 \int_0^x (x-t) y(t) dt + x \cdot \frac{\omega}{b} \int_b^b (b-t) y(t) dt \rightarrow \left(\int_0^x + \int_x^b \right)$$

Since

$$(x-t) = \frac{x}{b} (b-t) - \frac{t}{b} (b-x)$$

we find

$$y(x) = \omega^2 \int_0^x \frac{t}{b} (b-x) y(t) dt + \omega^2 \int_x^b \frac{x}{b} (b-t) y(t) dt$$

So, the kernel

$$K(x, t) = \begin{cases} \frac{t}{b} (b-x) & , t < x \\ \frac{x}{b} (b-t) & , t > x \end{cases}$$

$$y(x) = \omega^2 \int_0^b K(x, t) y(t) dt$$

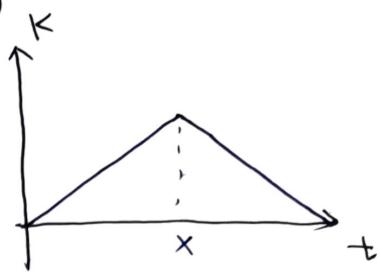
The kernel:

1. symmetric . $K(x, t) = K(t, x)$

2. continuous

3. $\frac{\partial K}{\partial t}$ is discontinuous

Green's function.



Initial condition \rightarrow Volterra

Boundary condition \rightarrow Fredholm .

differential \rightleftharpoons integral ,

Some special methods .

1. Integral transform.

$$f(x) = \int_{-\infty}^{+\infty} e^{ixt} \psi(t) dt \quad \text{Fourier}$$

$$\text{Solution } \psi(x) = \int_{-\infty}^{+\infty} e^{-ixt} f(t) dt$$

$$f(x) = \int_0^{+\infty} e^{xt} \psi(t) dt \quad \text{Laplace}$$

$$\psi(x) = \frac{1}{2\pi i} \int_{\beta-i\infty}^{\beta+i\infty} e^{xt} f(t) dt$$

$$f(x) = \int_0^x t^{x-1} \psi(t) dt$$

Mellin

$$\psi(x) = \frac{1}{2\pi i} \int_{\beta-i\infty}^{\beta+i\infty} x^{-t} f(t) dt$$

$$f(x) = \int_0^\infty t \psi(t) J_\nu(xt) dt$$

Hankel

$$\psi(x) = \int_0^\infty t f(it) J_\nu(xt) dt$$

Fourier Convolution theorem

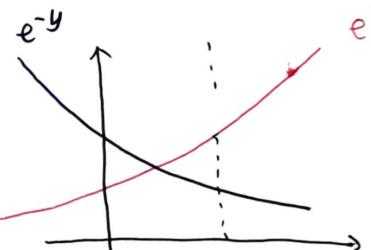
$$(f * g)(x) := \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} g(y) f(x-y) dy$$

3D:

$$(f * g)(\vec{r}) := \frac{1}{(2\pi)^{3/2}} \int g(\vec{r}') f(\vec{r} - \vec{r}') d^3 r'$$

Faltung : folding

$$\text{example: } f(y) = e^{-y}, \quad f(x-y) = e^{-(x-y)}$$

$$\begin{aligned} (f * g)^T(t) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} dx e^{itx} \left[\int_{-\infty}^{+\infty} dy g(y) f(x-y) \right] \\ &= \left[\frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} dy g(y) e^{ity} \right] \left[\frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} dx f(x-y) e^{it(x-y)} \right] \\ &= G(t) F(t) \end{aligned}$$


Inverse:

$$\frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} G(t) F(t) e^{-ixt} dt = (f * g)(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} g(y) f(x-y) dy$$

Laplace Convolution Theorem.

$$f_1(s) = \mathcal{L}\{F_1(t)\} \quad f_2(s) = \mathcal{L}\{F_2(t)\}$$

$$f_1(s)f_2(s) = \int_0^\infty e^{-sx} F_1(x) dx \int_0^\infty e^{-sy} F_2(y) dy$$

$$\text{let } t = x + y$$

$$\begin{aligned} f_1(s)f_2(s) &= \int_0^\infty e^{-st} dt \int_0^t F_1(t-y) F_2(y) dy \\ &= \mathcal{L}\left\{\int_0^t F_1(t-y) F_2(y) dy\right\} \\ &= \mathcal{L}\{F_1 * F_2\} \end{aligned}$$

$$F_1 * F_2 := \int_0^t F_1(t-z) F_2(z) dz$$

Inverse:

$$\mathcal{L}^{-1}\{f_1(s)f_2(s)\} = F_1 * F_2$$

Example. Fourier Transform Solution,

$$f(x) = \int_{-\infty}^{+\infty} k(x-t) \psi(t) dt$$

from fourier convolution theorem

$$f(x) = \int_{-\infty}^{+\infty} k(w) \Phi(w) e^{-iwx} dw$$

taking fourier transform.

$$k(w) \Phi(w) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} f(x) e^{iwx} dx = \frac{F(w)}{\sqrt{2\pi}}$$

$$\Phi(w) = \frac{1}{\sqrt{2\pi}} \frac{F(w)}{|k(w)|}$$

$$\psi(x) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} \frac{F(w)}{|k(w)|} e^{-iwx} dw$$

example: Generalized Abel Equation

$$f(x) = \int_0^x \frac{\psi(t)}{(x-t)^\alpha} dt, \quad 0 < \alpha < 1$$

Taking the Laplace transform.

$$\mathcal{L}\{f(x)\} = \mathcal{L}\left\{\int_0^x \frac{\psi(t)}{(x-t)^\alpha} dt\right\} = \mathcal{L}\{x^{-\alpha}\} \mathcal{L}\{\psi(x)\}$$

$$\mathcal{L}\{\psi(x)\} = \frac{s^{1-\alpha}}{\Gamma(1-\alpha)} \cdot \mathcal{L}\{f(x)\}$$

$$\frac{1}{s} \mathcal{L}\{\psi(x)\} = \frac{s^{-\alpha}}{\Gamma(1-\alpha)} \mathcal{L}\{f(x)\} = \frac{\mathcal{L}\{x^{\alpha-1}\} \mathcal{L}\{f(x)\}}{\Gamma(\alpha) \Gamma(1-\alpha)}$$

$$\downarrow$$

$$\mathcal{L}\left\{\int_0^x \psi(t) dt\right\} = \frac{\sin \pi \alpha}{\pi} \mathcal{L}\left\{\int_0^x \frac{f(t)}{(x-t)^{1-\alpha}} dt\right\}$$

$$\int_0^x \psi(t) dt = \frac{\sin \pi \alpha}{\pi} \int_0^x \frac{f(t)}{(x-t)^{1-\alpha}} dt$$

$$\psi(t) = \frac{\sin \pi \alpha}{\pi} \frac{d}{dx} \int_0^x \frac{f(t)}{(x-t)^{1-\alpha}} dt$$

Generating Function Method.

$$f(x) = \int_{-1}^1 \frac{\psi(t)}{\sqrt{1-2xt+x^2}} dt, \quad -1 \leq x \leq 1 \Rightarrow$$

Legendre polynomials.

$$\frac{1}{\sqrt{1-2xt+x^2}} = \sum_{n=0}^{\infty} P_n(t) x^n$$

$$\psi(t) = \sum_{m=0}^{\infty} a_m P_m(t)$$

$$\therefore f(x) = \sum_{mn} a_m x^n \int_{-1}^1 P_n(t) P_m(t) dt$$

$$= \sum_{mn} a_m x^n \frac{2 \delta_{mn}}{2n+1} = \sum_{n=0}^{\infty} \frac{2 a_n}{2n+1} x^n$$

orthonormal

Maclaurin expansion for $f(x)$

$$f(x) = \sum_{n=0}^{\infty} \frac{f^{(n)}(0)}{n!}$$

So, equate them.

$$\frac{f^{(n)}(0)}{n!} = \frac{2a_n}{2n+1}$$

$$\begin{aligned}\psi(t) &= \sum_{m=0}^{\infty} a_m P_m(t) \\ &= \sum_{n=0}^{\infty} \frac{2n+1}{2} \frac{f^{(n)}(0)}{n!} P_n(t)\end{aligned}$$

Separable Kernel

$$K(x, t) = \sum_{j=1}^n M_j(x) N_j(t), n < +\infty$$

degenerate

Polynomials and some elementary transdental function

as.

$$\begin{aligned}K(x, t) &= \cos(t - x) \\ &= \cos t \cos x + \sin t \sin x\end{aligned}$$

For a fredholm equation . (2)

$$\psi(x) = f(x) + \lambda \sum_{j=1}^n M_j(x) \underbrace{\int_a^b N_j(t) \psi(t) dt}_{c_j}$$

So,

$$\psi(x) = f(x) + \lambda \sum_{j=1}^n c_j M_j(x)$$

multiply $N_i(x)$ and integrate

$$\underbrace{\int_a^b \varphi_i(x) N_i(x) dx}_{c_i} = \underbrace{\int_a^b f(x) N_i(x) dx}_{b_i} + \lambda \sum_{j=1}^n c_j \underbrace{\int_a^b M_j(x) N_i(x) dx}_{a_{ij}}$$

$$c_i = b_i + \lambda \sum_j a_{ij} c_j$$

$$\vec{c} = \vec{b} + \lambda A \vec{c}$$

$$\vec{b} = (\mathbb{1} - \lambda A) \vec{c}$$

$$\vec{c} = (\mathbb{1} - \lambda A)^{-1} \vec{b}$$

If $f(x) = 0$, then $\vec{b} = \vec{0}$

$$|\mathbb{1} - \lambda A| = 0$$

example

$$\varphi_i(x) = \lambda \int_1^x (t+x) \varphi(t) dt$$

$$t+x = 1 \cdot t + x \cdot 1 = M_1(x) N_1(t) + M_2(x) N_2(t)$$

$$a_{ij} = \int_1^x N_i(x) M_j(x) dx$$

$$\text{So, } \begin{vmatrix} 1 & -\frac{2\lambda}{3} \\ -2\lambda & 1 \end{vmatrix} = 0$$

$$\lambda = \pm \frac{\sqrt{3}}{2}$$

$$\vec{b} = (\mathbb{1} - \lambda A) \vec{c} = \vec{0} \Rightarrow c_1 + \frac{c_2}{\sqrt{3}} = 0$$

If let $c_1 = 1$

$$\varphi_1(x) = \frac{\sqrt{3}}{2} (1 + \sqrt{3}x) \quad \lambda = \frac{\sqrt{3}}{2}$$

$$\varphi_2(x) = -\frac{\sqrt{3}}{2} (1 - \sqrt{3}x) \quad \lambda = -\frac{\sqrt{3}}{2}$$

21.3. Neumann Series

$$\varphi(x) = f(x) + \lambda \int_a^b K(x, t) \varphi(t) dt, \quad f(x) \neq 0$$

Suppose

$$\varphi_0(x) \approx f(x)$$

$$\varphi_1(x) = f(x) + \lambda \int_a^b K(x, t) f(t) dt$$

$$\begin{aligned} \varphi_2(x) &= f(x) + \lambda \int_a^b K(x, t_1) f(t_1) dt_1 \\ &\quad + \lambda^2 \int_a^b \int_a^b K(x, t_1) K(t_1, t_2) f(t_2) dt_2 dt_1 \end{aligned}$$

$$\varphi_n(x) = \sum_{i=0}^n \lambda^i u_i(x)$$

where

$$u_0(x) = f(x)$$

$$u_1(x) = \int_a^b K(x, t_1) f(t_1) dt_1,$$

$$u_n(x) = \underbrace{\int_a^b \dots \int_a^b}_{n} \left[\prod_{i=1}^n K_i \right] f(t_n) dt_n \dots dt_1$$

We expect

$$\varphi(x) = \lim_{n \rightarrow \infty} \varphi_n(x) = \lim_{n \rightarrow \infty} \sum_{i=0}^n \lambda^i u_i(x)$$

$$|\lambda^n u_n(x)| \leq |\lambda^n| \cdot \|f\|_{\max} \cdot \|K\|_{\max}^n \cdot |b-a|^n$$

So, if $|\lambda| \|K\|_{\max} |b-a| < 1$, converge.

operator form

$$\varphi = \lambda \hat{K} \varphi + f$$

$$\hat{K} = \int_a^b K(x, t) [] dt$$

$$\Rightarrow \varphi = (\mathbb{1} - \lambda \hat{K})^{-1} f$$

Binomial expansion

example.

$$\varphi_1(x) = x + \frac{1}{2} \int_{-1}^1 (t-x) \varphi(t) dt$$

$$\varphi_1(x) = x$$

Then

$$\varphi_2(x) = x + \frac{1}{2} \int_{-1}^1 (t-x)t dt = x + \frac{1}{3}$$

$$\varphi_3(x) = x + \frac{1}{3} - \frac{x}{3}$$

$$\varphi_4(x) = x + \frac{1}{3} - \frac{x}{3} - \frac{1}{3^2}$$

:

$$\varphi(x) = \frac{3}{4}x + \frac{1}{4} .$$