10.5:

We can derive that

$$\frac{\partial S(q,\alpha,t)}{\partial t} = -\frac{1}{2}m\omega^2(\alpha^2 + q^2 - 2\alpha q\cos\omega t)\csc^2\omega t$$

and

$$\frac{\partial S(q,\alpha,t)}{\partial q} = m\omega(q\cos\omega t - a)\csc\omega t.$$

It is obvious that we can write the Hamiltonian as,

$$H\left(q,\frac{\partial S}{\partial q},t\right) = \frac{1}{2}m\omega^{2}\left[(q\cos\omega t - a)^{2}\csc^{2}\omega t + q^{2}\right]$$

Therefore, the principal function satisfies,

$$\frac{\partial S(q,\alpha,t)}{\partial t} + H\left(q,\frac{\partial S}{\partial q},t\right) = 0,$$

which is a solution of the Hamiltonian-Jacobi. And clearly, by the Hamilton's equations,

$$\dot{p} = -\frac{\partial H}{\partial q} = -m\omega^2 q,$$
$$\dot{q} = \frac{\partial H}{\partial p} = \frac{p}{m}.$$

It leads that

$$\ddot{q} + \omega^2 q = 0.$$

Thus, this function generates the solution to the motion of the harmonic oscillator.

10.11:

We have

$$\dot{z} = -\frac{4\pi \dot{x}}{\lambda} A \cos \frac{2\pi x}{\lambda} \sin \frac{2\pi x}{\lambda}.$$

Thus,

$$L = \frac{1}{2}m\dot{x}^2 \left(1 + \frac{16\pi^2 A^2}{\lambda^2}\cos^2\frac{2\pi x}{\lambda}\sin^2\frac{2\pi x}{\lambda}\right) - mg\cos^2\frac{2\pi x}{\lambda}.$$

And as a result, we have

$$p = \frac{\partial L}{\partial \dot{x}} = m\dot{x} \left(1 + \frac{16\pi^2 A^2}{\lambda^2} \cos^2 \frac{2\pi x}{\lambda} \sin^2 \frac{2\pi x}{\lambda} \right).$$

And

$$H = p\dot{x} - L = \frac{p^2}{2m\left(1 + \frac{16\pi^2 A^2}{\lambda^2}\cos^2\frac{2\pi x}{\lambda}\sin^2\frac{2\pi x}{\lambda}\right)} + mgA\cos^2\frac{2\pi x}{\lambda}.$$

And naturally, we have

$$\frac{dx}{dt} = \frac{\partial H}{\partial p} = \frac{p}{m\left(1 + \frac{4\pi^2 A^2}{\lambda^2}\sin^2\frac{4\pi x}{\lambda}\right)}.$$

$$\frac{dp}{dt} = -\frac{\partial H}{\partial x} = -\frac{p^2}{m\left(1 + \frac{4\pi^2 A^2}{\lambda^2}\sin^2\frac{4\pi x}{\lambda}\right)^2} \frac{4\pi^2 A^2 \sin\frac{4\pi x}{\lambda}\cos\frac{4\pi x}{\lambda}}{\lambda^2} - \frac{4\pi}{\lambda} mgA\cos\frac{2\pi x}{\lambda}\sin\frac{2\pi x}{\lambda}.$$

This is the evolution of the system in the phase space.

10.14:

The Hamiltonian is

$$H\left(x,\frac{\partial S}{\partial x}\right) = \frac{1}{2m}\left(\frac{\partial S}{\partial x}\right) - \frac{k}{|x|}.$$

The Hamilton-Jacobi equation yields

$$\frac{1}{2m}\left(\frac{\partial S}{\partial x}\right) - \frac{k}{|x|} + \frac{\partial S}{\partial t} = 0.$$

Because the energy is conserved, we can take the principal function as the form

$$S = W(x, E) - Et + C.$$

Providing the energy is negative,

$$\frac{\partial W}{\partial x} = \sqrt{2m} \left[\frac{k}{|x|} - |E| \right]^{1/2}.$$

The turning point is at $x = \pm k/|E|$, and the action variable is,

$$J = \oint p dq = 4\sqrt{2m} \int_0^{k/|E|} \left[\frac{k}{x} - E\right]^{1/2} dx = 2\pi k \sqrt{\frac{2m}{|E|}}.$$

Then we can express the Hamiltonian as

$$H = E = -\frac{8m\pi^2k^2}{J^2}.$$

The period is

$$T = 1/\nu = \left(\frac{\partial H}{\partial J}\right)^{-1} = \left(\frac{16m\pi^2 k^2}{J^3}\right)^{-1} = \frac{\pi k\sqrt{2m}}{|E|^{3/2}}.$$

10.17:

The Hamiltonian of projectile is given by

$$H = \frac{p_x^2}{2m} + \frac{p_y^2}{2m} + mgy.$$

The Hamiltonian-Jacobi equation reads

$$\frac{1}{2m}\left(\frac{\partial S}{\partial x}\right)^2 + \frac{1}{2m}\left(\frac{\partial S}{\partial y}\right)^2 + mgy + \frac{\partial S}{\partial t} = 0.$$

Since the Hamiltonian does not depend on time explicitly, we can write the principal function as

$$S = -Et + W(x, y) + C.$$

Therefore, it turns out

$$\frac{1}{2m} \left(\frac{\partial W}{\partial x}\right)^2 + \frac{1}{2m} \left(\frac{\partial W}{\partial y}\right)^2 + mgy = E$$

Since the Hamiltonian can be decomposed as

$$H = f_1\left(x, \frac{\partial W}{\partial x}\right) + f_2\left(y, \frac{\partial W}{\partial y}\right),$$

we can also decompose \boldsymbol{W} as

$$W = W_1(x) + W_2(y).$$

We have

$$\frac{1}{2m} \left(\frac{dW_1}{dx}\right)^2 = \alpha_1 \implies W_1 = \pm \sqrt{2\alpha_1 m x},$$
$$\frac{1}{2m} \left(\frac{dW_2}{dy}\right)^2 + mgy = \alpha_2 \implies W_2 = \pm \frac{2\sqrt{2}(\alpha_2 - mgy)^2}{3g\sqrt{m(\alpha_2 - mgy)^2}},$$
$$\alpha_1 + \alpha_2 = E.$$

Therefore, the principal function is

$$S = -Et \pm \sqrt{2\alpha_1 m} x \pm \frac{2\sqrt{2}(\alpha_2 - mgy)^2}{3g\sqrt{m(\alpha_2 - mgy)}} + C.$$

We can derive

$$\beta_1 = \frac{\partial S}{\partial \alpha_1} = -t \pm \sqrt{\frac{m}{2\alpha_1}} x \implies x = \pm \sqrt{\frac{2\alpha_1}{m}} (\beta_1 + t),$$

$$\beta_2 = \frac{\partial S}{\partial \alpha_2} = -t \pm \frac{\sqrt{2m}}{mg} \sqrt{\alpha_2 - mgy} \implies y = \frac{\alpha_2}{mg} - \frac{g}{2} (\beta_2 + t)^2.$$

Putting initial conditions, we get

$$x = v_0 t \cos \alpha,$$

$$y = v_0 t \sin \alpha - \frac{1}{2}gt^2.$$

10.26:

The Lagrangian is

$$L = \frac{1}{2}I_1(\dot{\theta}^2 + \psi \sin^2 \theta) + \frac{1}{2}I_3(\dot{\phi} + \psi \cos \theta)^2 - Mgl\cos\theta.$$

The conjugate momenta are

$$p_{\theta} = I_1 \dot{\theta},$$

$$p_{\psi} = I_3 (\dot{\psi} + \dot{\phi} \cos \theta),$$

$$p_{\phi} = I_3 \dot{\psi} \cos \theta + \dot{\phi} (I_1 \sin^2 \theta + I_3 \cos^2 \theta).$$

After that, we have

$$E = \frac{1}{2}I_3\omega_3^2 + \frac{I_1}{2}\frac{(b-a\cos\theta)^2}{\sin^2\theta} + Mgl\cos\theta.$$
$$I_1a = p_{\phi},$$
$$I_2b = p_{\psi}.$$

We assume that

$$\alpha = \frac{I_1}{2} \frac{(b - a\cos\theta)^2}{\sin^2\theta} + Mgl\cos\theta.$$

After separation of variables, we have

$$W = W(\theta) + W(\psi) + W(\phi).$$

Where

Where

$$W(\psi) = I_1 a \psi,$$

$$W(\phi) = I_1 b \phi.$$

And

$$W(\theta) = \int d\theta \sqrt{2I_1\alpha - I_1^2 \frac{(b - a\cos\theta)^2}{\sin^2\theta} - 2I_1 Mgl\cos\theta}.$$

And

$$t + \beta_1 = \frac{\partial W}{\partial \alpha},$$

$$\beta_2 = \frac{\partial W}{\partial a},$$

$$\beta_3 = \frac{\partial W}{\partial b}.$$

And after rearranging parameters, we have

$$\dot{u}^2 = (1 - u^2) (\alpha - \beta u) - (b - au)^2.$$

Where

$$\alpha = \frac{2E - I_3 \omega_3^2}{I_1},$$
$$\beta = \frac{2Mgl}{I_1}.$$