

Maxwell Equations. and Proca Equations.

1) Maxwell equations.

$$\begin{cases} \nabla \cdot \mathbf{B} = 0 & \textcircled{1} \\ \nabla \times \mathbf{B} = \mathbf{j} + \frac{\partial \mathbf{E}}{\partial t} & \textcircled{2} \\ \nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t} & \textcircled{3} \\ \nabla \cdot \mathbf{E} = \rho & \textcircled{4} \end{cases}$$

Vector form $\xrightarrow{\text{SR}}$ Covariant form.

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Lorentz trans.

introduce : Magnetic vector potential : \mathbf{A}

$$\begin{cases} \mathbf{B} = \nabla \times \mathbf{A} \\ \mathbf{E} = -\frac{\partial \mathbf{A}}{\partial t} - \nabla \phi \end{cases}$$

4-Vector $A^\mu = (\phi, \vec{A})$

$$\begin{cases} E^i = \partial^0 A^i - \partial^i A^0 \\ B^i = \epsilon^{ijk} \partial^j A^k \end{cases}$$

EMF Tensor: $F^{\mu\nu} = -F^{\nu\mu} = \partial^\mu A^\nu - \partial^\nu A^\mu =$

$$\begin{pmatrix} 0 & -E^1 & -E^2 & -E^3 \\ E^1 & 0 & -B^3 & B^2 \\ E^2 & B^3 & 0 & -B^1 \\ E^3 & -B^2 & B^1 & 0 \end{pmatrix}$$

1). Anti-sys.

2). Traceless

under Lorentz transformation:

$$F^{\mu\nu} \rightarrow \Lambda^\mu_\alpha \Lambda^\nu_\beta F^{\alpha\beta}$$

2) Gauge

$$\begin{aligned} \mathbf{A} &\rightarrow \mathbf{A} - \nabla \chi \\ \phi &\rightarrow \phi + \frac{\partial \chi}{\partial t} \end{aligned} \Rightarrow A^\mu \rightarrow A^\mu + \partial^\mu \chi$$

$$F^{\mu\nu} = F^{\mu\nu} + (\partial^\mu \partial^\nu - \partial^\nu \partial^\mu) \chi = F^{\mu\nu}$$

$$\square A^\mu = j^\mu + \partial^\nu (\partial_\nu A^\mu)$$

↓
0

Lorentz gauge

$$\partial_\mu A^\mu = \frac{\partial \phi}{\partial t} + \nabla \cdot \mathbf{A} = 0$$

$$\square A^\mu = j^\mu \quad 0$$

For vac

$$\begin{cases} \frac{\partial \phi}{\partial t^2} - \nabla^2 \phi = \rho & (\text{Liénard-Wiechert potential.}) \\ \frac{\partial A}{\partial t^2} - \nabla^2 A = \vec{j} \end{cases}$$

$$\square A^\mu = 0$$

Back to Maxwell equation:

$$\partial_\mu F^{\mu\nu} = j^\nu \quad \star$$

$$\Rightarrow \begin{cases} \nu=0 : \nabla \cdot E = \rho & \textcircled{2} \\ \nu=1,2,3 : \nabla \times B = \vec{j} + \frac{\partial E}{\partial t} & \textcircled{4} \end{cases}$$

$$F = -\frac{1}{2} F_{\mu\nu} dx^\mu \wedge dx^\nu$$

$$\star F^{\mu\nu} = G^{\mu\nu} = \frac{1}{2} \epsilon^{\mu\nu\rho\sigma} F_{\rho\sigma} = \begin{pmatrix} 0 & -B^1 & -B^2 & -B^3 \\ B^1 & 0 & E^3 & -E^2 \\ B^2 & -E^3 & 0 & E^1 \\ B^3 & E^2 & -E^1 & 0 \end{pmatrix}$$

4-D Levi-Civita

$$\partial_\mu G^{\mu\nu} = 0 \Rightarrow \textcircled{1} \textcircled{3}$$

Maxwell equation \rightarrow Tensor form.

$$\begin{cases} \partial_\mu F^{\mu\nu} = j^\nu \\ \partial_\mu G^{\mu\nu} = 0 \end{cases}$$

\Rightarrow Proca Equations: For massive spin-1 particles.

$$F^{\mu\nu} = \partial^\mu A^\nu - \partial^\nu A^\mu$$

$$\partial_\mu F^{\mu\nu} + m^2 A^\nu = 0$$

\uparrow
mass

$$m^2 \partial_\mu A^\nu = 0 \quad (\square + m^2) A^\mu = 0 \quad (\partial_\mu A^\mu = 0)$$

Proca Equation $\begin{cases} (\square + m^2) A^\mu = 0 \\ \partial_\mu F^{\mu\nu} + m^2 A^\nu = 0 \end{cases}$

2.9. Differential Form.

Intro: $\int I_1 = \int_C F_x dx + F_y dy + F_z dz = \int_C F \cdot dr = \mathcal{R}$
 $\int I_2 = \int_S (G_x dy \wedge dz + G_y dz \wedge dx + G_z dx \wedge dy)$
 $= \int_S G \cdot ds = \mathcal{R}$ \downarrow
 orientation

$C \leftrightarrow S$
 \uparrow dual \uparrow

1. $C_n \rightarrow n$ -D volume
2. Boundary Operator: $\partial C_n \rightarrow C_{n-1}$ $\partial C_n \rightarrow C_{n-1}$
3. $\partial \mathbb{Z}_n = 0$ \mathbb{Z}_n . no boundary
4. $B_n = \partial C_{n+1} \Rightarrow \partial B_n = 0$
 $\left. \begin{array}{l} \partial^2 = 0 \\ \mathbb{Z}_n \end{array} \right\}$ for Euclidean $\mathbb{Z}_n = B_n$, $\mathbb{Z}_n \supset B_n$

5. w_0 0-form \rightarrow function \hookleftarrow 3-D
 w_1 1-form $\rightarrow A dx + B dy + C dz$
 w_2 2-form $\rightarrow f dx \wedge dy + g dy \wedge dz + h dz \wedge dx$
 w_3 3-form $\rightarrow F dx \wedge dy \wedge dz$

$dx \wedge dy = -dy \wedge dx$, $dx \wedge dx = 0$

b. Exterior derivative operator:

$dw_n = w_{n+1}$

3-D

$d(A dx + B dy + C dz) = \left(\frac{\partial B}{\partial x} - \frac{\partial A}{\partial y}\right) dx \wedge dy + \left(\frac{\partial C}{\partial y} - \frac{\partial B}{\partial z}\right) dy \wedge dz + \left(\frac{\partial A}{\partial z} - \frac{\partial C}{\partial x}\right) dz \wedge dx$

$d(A, B, C) \rightarrow \nabla \times (A, B, C)$

$dw_2 = \left(\frac{\partial f}{\partial z} + \frac{\partial g}{\partial x} + \frac{\partial h}{\partial y}\right) dx \wedge dy \wedge dz$

$\nabla \cdot (g, h, f)$

$d^2 = 0$ ($\nabla \cdot (\nabla \times \vec{v}) = 0$)

$$\int d\omega_n = 0 \Rightarrow \omega_n \text{ is closed}$$

$$\int \omega_n = dW_{n-1} \Rightarrow \omega_n \text{ is exact}$$

Stokes : $\int_{\partial C} \omega = \int_C d\omega$

4) Hodge dual star operator

For k -d Vector space (inner product).

$$\langle \alpha, \beta \rangle = \det (\langle \alpha_i, \beta_j \rangle_{i,j=1}^k) \quad \alpha, \beta \in \Lambda^k V \quad (0 \leq k \leq n)$$

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$\{e_1, e_2, \dots, e_n\}$ of V as:

$$\omega := e_1 \wedge e_2 \wedge \dots \wedge e_n$$


Hodge : $d \wedge (*\beta) = \langle \alpha, \beta \rangle \omega$

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$$k\text{-D} \rightarrow (n-k)\text{-D}$$

$$\beta \rightarrow *\beta$$

$$*(\omega_1 \wedge \dots \wedge \omega_k) = \omega_1 \wedge \dots \wedge \omega_{n-k}$$

Example :

2-D

$$\left\{ \begin{aligned} *1 &= dx \wedge dy \\ *dx &= dy \\ *dy &= -dx \\ *(dx \wedge dy) &= 1 \end{aligned} \right.$$

4-D $*dx^\mu = \eta^{\mu\nu} \epsilon_{\lambda\nu\rho\sigma} \frac{1}{3!} dx^\lambda \wedge dx^\rho \wedge dx^\sigma$

$$*(dx^\mu \wedge dx^\nu) = \eta^{\mu\kappa} \eta^{\nu\lambda} \epsilon_{\lambda\rho\sigma\kappa} \frac{1}{2!} dx^\rho \wedge dx^\sigma$$



$$F = -\frac{1}{2} F_{\mu\nu} dx^\mu \wedge dx^\nu$$

$$*F = -\frac{1}{2} F_{\mu\nu} *(dx^\mu \wedge dx^\nu)$$

$$\begin{cases} d*F = J \\ dF = 0 \end{cases}$$

$$J = (j_x dy \wedge dz + j_y dz \wedge dx + j_z dx \wedge dy) \wedge dt - \rho dx \wedge dy \wedge dz$$