

# Electro-Weak Theory

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$SU(2)_L \times U(1)_Y$  symmetry.  
neutral current ...?

1.  $SU(2)_L \times U(1)_Y$  symmetry.

leptons:  $\begin{pmatrix} e \\ \nu_e \end{pmatrix} \begin{pmatrix} \mu \\ \nu_\mu \end{pmatrix} \begin{pmatrix} \tau \\ \nu_\tau \end{pmatrix}$  quarks:  $\begin{pmatrix} u \\ d \end{pmatrix} \begin{pmatrix} c \\ s \end{pmatrix} \begin{pmatrix} t \\ b \end{pmatrix}$

$SU(2)_L \rightarrow SU(2)_L$  doublet  $L = \begin{pmatrix} \nu_e \\ e \end{pmatrix}_L$   $Q = \begin{pmatrix} c \\ s \end{pmatrix}_L$  isospin: 同位旋.  
 $SU(2)_R$  singlet  $e_R \ \mu_R \ \tau_R \ u_R \ d_R$

$\mathcal{L}_{Fermions} = \underbrace{\bar{L} i \not{\partial} L}_{SU(2)_L} + \underbrace{\bar{Q} i \not{\partial} Q + \bar{e}_R i \not{\partial} e_R + \bar{\mu}_R i \not{\partial} \mu_R + \bar{\tau}_R i \not{\partial} \tau_R}_{SU(2)_R} + \dots$  Gauge problem!

gauge transformation:  $\psi_L = L/R$ .  $\psi_L$  from  $\frac{1}{2}(1-\gamma_5)\psi$ .  
 $\begin{cases} \psi_L \rightarrow e^{i\alpha^a \tau^a / 2} e^{iY\psi/2} \psi_L & e^* e^Y \neq e^{*+Y} \text{ (Here is OK!)} \\ \psi_L \rightarrow e^{i[\alpha^a \tau^a + Y\psi/2]} \psi_L \\ \psi_R \rightarrow e^{iY\psi/2} \psi_R \end{cases}$   
 $Y_L, Y_R$ : hypercharge (number)  
 $Y_L = Y_R = Y$ : 盖尔曼-西岛关系:  $Y/2 = Q - I_3$ .

for  $\psi_L$ :  $I_3 = \begin{cases} \frac{1}{2} & \text{v. u.c.t messy!} \\ -\frac{1}{2} & \text{e. p. v. d. s. b.} \end{cases}$

Is  $U(1)_Y$  same as  $U(1)_Q$ ? The answer is not.

$Q \rightarrow$  doublet:  $\begin{cases} \begin{pmatrix} 0 & -1 \\ \frac{2}{3} & -\frac{1}{3} \end{pmatrix} & \text{for leptons} \\ \begin{pmatrix} 2/3 & -1/3 \end{pmatrix} & \text{for quarks} \end{cases}$   $SU(2)_L$ : generators are  $\tau^a/2$

they don't commute?  $\Rightarrow U(1)_Q \times SU(2)_L$  is bad!

$\mathcal{L}_{Gauge} + \mathcal{L}_{Fermions} = -\frac{1}{4} F_{\mu\nu}^a F^{\mu\nu a} - \frac{1}{4} B_{\mu\nu} B_{\mu\nu} + \mathcal{L}'_{Fermions}$ . prime means  $\not{\partial} \rightarrow \not{D}$ .  
 $\Rightarrow \mathcal{L}'_{Fermions} = \bar{e}_R i \not{\partial} e_R + \bar{L} i \not{\partial} L$  interactions  
 $\begin{cases} D_\mu^L = \partial_\mu - ig \frac{\tau^a}{2} W_\mu^a - ig \frac{Y}{2} B_\mu \\ D_\mu^R = \partial_\mu - ig \frac{Y}{2} B_\mu \end{cases}$   $W_\mu^a$  for  $F_{\mu\nu}^a$   
 $B_\mu$  for  $B_{\mu\nu}$

Define:  $W_\mu^\pm = \frac{1}{\sqrt{2}}(W_\mu^1 \mp iW_\mu^2)$   $W_\mu^3 \rightarrow A_\mu$   $B_\mu \rightarrow Z_\mu$

charged weak intermediate bosons photon special field: coupling to weak neutral currents.

How to get their masses? SSB: Higgs Mechanism.

$SU(2)_L \times U(1)_Y \rightarrow U(1)_Q$

$\mathcal{L} = \mathcal{L}_{Gauge} + \mathcal{L}_{Fermions} + \mathcal{L}_S$   $\mathcal{L}_S = D_\mu \phi^\dagger D^\mu \phi - V(\phi)$ .  $\phi$ : should be  $SU(2)_L$  doublet.  
 $\Rightarrow D_\mu = (\partial_\mu - ig \frac{\tau^a}{2} W_\mu^a - ig \frac{Y}{2} B_\mu)$   $\phi \rightarrow e^{i\alpha^a \tau^a / 2} e^{iY\phi/2} \phi$

$\rho(x)$ .  $\alpha^i(x) \rightarrow$  coordinate dependent.

SSB calculation: (1). Vacuum expectation value:  $\langle \phi \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v \end{pmatrix}$

(2) expand near the vacuum:

$\Delta \mathcal{L} = \frac{1}{2} \left[ \frac{g^2}{4} (W_\mu^1)^2 + g^2 (W_\mu^2)^2 + (-gW_\mu^3 + g' B_\mu)^2 \right]$

recall that:  $W_\mu^\pm = \frac{1}{\sqrt{2}}(W_\mu^1 \mp iW_\mu^2) \Rightarrow m_W = \frac{g v}{2}$  (W boson)

$g, g'$ :  $\cos \theta_w = \frac{g}{\sqrt{g^2 + g'^2}}$   $\sin \theta_w = \frac{g'}{\sqrt{g^2 + g'^2}}$   $\theta_w$ : Weinberg angle

$\left. \begin{aligned} B_\mu &= \cos \theta_w A_\mu - \sin \theta_w Z_\mu \\ W_\mu^3 &= \sin \theta_w A_\mu + \cos \theta_w Z_\mu \end{aligned} \right\} \sim \frac{g^2}{\cos^2 \theta_w} Z_\mu Z_\mu \Rightarrow m_Z = \frac{m_W}{\cos \theta_w}$  (Z boson).

from 4-fermions theory:  $G_F \approx 1.166 \times 10^{-5} \text{ GeV}^{-2}$   $\frac{G_F}{\sqrt{2}} = \frac{1}{2v^2} \Rightarrow v \approx 246 \text{ GeV}$ .

from experiment:  $\sin^2 \theta_w \approx 0.23 \Rightarrow \begin{cases} m_W \approx 80 \text{ GeV} \\ m_Z \approx 90 \text{ GeV} \end{cases}$   $m_A = 0$  massless.

Why EM-field?  $\tau^3 \begin{pmatrix} 0 \\ v \end{pmatrix} \neq 0$   $Y \begin{pmatrix} 0 \\ v \end{pmatrix} \neq 0$

$Q = \left( \frac{\tau^3}{2} + \frac{Y}{2} \right) = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$   $Q \begin{pmatrix} 0 \\ v \end{pmatrix} = 0$ . EM charge satisfy  $U(1)_Q$  symmetry.

Gauge sector. Higgs sector. Fermions sector